

Inverse source problems of local, nonlocal and nonlinear equations¹

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6th Young Scholar Symposium, March 25-26, 2023.

The Chinese University of Hong Kong



科技部年輕學者養成計畫
TAIWAN MOST Young Scholar Fellowship

¹This talk is based on recent different joint works with [Yavar Kian](#), [Tony Liimatainen](#) and [Hongyu Liu](#).

Outline

- 1 Inverse source problems for nonlinear equations
- 2 The fractional Calderón problem
- 3 Inverse source problems for nonlocal equations

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Inverse source problems for elliptic equations

Consider

$$\Delta u = F \text{ in } \Omega, \quad u = f \text{ on } \partial\Omega. \quad (1.1)$$

An inverse source problem is to determine the source F by using the DN map. However, it can be seen that it is impossible to determine F due to the **gauge invariance**.

- **Gauge invariance.** Consider $\phi \in C_c^2(\Omega)$, and let $\tilde{u} := u + \phi$, then $\Delta \tilde{u} = F + \Delta \phi$. In general, $F + \Delta \phi \neq F$, but the Cauchy data of \tilde{u} and u are the same, i.e.,

$$\{\tilde{u}|_{\partial\Omega}, \partial_\nu \tilde{u}|_{\partial\Omega}\} = \{u|_{\partial\Omega}, \partial_\nu u|_{\partial\Omega}\}.$$

Thus, it is natural to study similar questions for inverse source problems for **nonlinear equations**.

Inverse source problems for nonlinear elliptic equations

Theorem (Liimatainen-L., gauge invariance 2022)

Let $a_j(x, z) = a_j^{(1)}(x)z + a_j^{(2)}(x)z^2$, where $a_j^{(1)}, a_j^{(2)} \in C^\alpha(\bar{\Omega})$ for some $0 < \alpha < 1$, for $j = 1, 2$. Consider

$$\begin{cases} \Delta u_j + a_j(x, u_j) = F_j & \text{in } \Omega, \\ u_j = f & \text{on } \partial\Omega, \end{cases} \quad (1.2)$$

and let Λ_{a_j, F_j} to be the corresponding DN map of (1.2) for $j = 1, 2$. Suppose that

$$\Lambda_{a_1, F_1}(f) = \Lambda_{a_2, F_2}(f) \text{ for any } f \in \mathcal{N}.$$

Then there exists $\psi \in C^{2,\alpha}(\bar{\Omega})$ with $\psi|_{\partial\Omega} = \partial_\nu \psi|_{\partial\Omega} = 0$ in Ω such that

$$\begin{cases} a_1^{(2)} = a_2^{(2)} =: a^{(2)}, & a_1^{(1)} = a_2^{(1)} + 2a^{(2)}\psi, \\ F_1 = F_2 - \Delta\psi - a_1^{(2)}\psi - a^{(2)}\psi^2. \end{cases} \quad (1.3)$$

In the nonlinear counterpart, we can further get **uniqueness** result for both coefficients and sources.

Corollary (Liimatainen-L., gauge breaking 2022)

For the quadratic case, assume additionally that

$$a_1^{(1)} = a_2^{(1)} \text{ in } \Omega$$

and

$$a_1^{(2)}(x) \neq 0 \text{ or } a_2^{(2)}(x) \neq 0 \text{ at any } x \in \Omega.$$

Then also

$$F_1 = F_2 \text{ and } a_1^{(2)} = a_2^{(2)} \text{ in } \Omega.$$

The proofs of both theorem and corollary are based on the **higher order linearization**.

The proof

1. Initiation.

We apply the **higher order linearization** method to the equation

$$\begin{cases} \Delta u_j + a_j^{(1)} u_j + a_j^{(2)} u_j^2 = F_j & \text{in } \Omega, \\ u_j = f_0 + \varepsilon_1 f_1 + \varepsilon_2 f_2 & \text{on } \partial\Omega. \end{cases} \quad (1.4)$$

We denote $\varepsilon = (\varepsilon_1, \varepsilon_2)$, which especially means that $\varepsilon = 0$ is equivalent to $\varepsilon_1 = \varepsilon_2 = 0$. Below the index $j = 1, 2$ corresponds to the different sets of coefficients, and an index $\ell = 1, 2$ to ε_ℓ parameters. Let us denote by $u_j^{(0)}$ the solution to

$$\begin{cases} \Delta u_j^{(0)} + a_j^{(1)} u_j^{(0)} + a_j^{(2)} \left(u_j^{(0)}\right)^2 = F & \text{in } \Omega, \\ u_j^{(0)} = f_0 & \text{on } \partial\Omega. \end{cases} \quad (1.5)$$



2. First linearization.

Differentiate (1.4) with respect to ε_ℓ , for $\ell = 1, 2$. We obtain

$$\begin{cases} \left(\Delta + a_j^{(1)} + 2a_j^{(2)} u_j^{(0)} \right) v_j^{(\ell)} = 0 & \text{in } \Omega, \\ v_j^{(\ell)} = f_\ell & \text{on } \partial\Omega, \end{cases} \quad (1.6)$$

where

$$v_j^{(\ell)} = \partial_{\varepsilon_\ell} |_{\varepsilon=0} u_j,$$

for $j, \ell = 1, 2$. The global uniqueness implies

$$Q := a_1^{(1)} + 2a_1^{(2)} u_1^{(0)} = a_2^{(1)} + 2a_2^{(2)} u_2^{(0)} \text{ in } \Omega. \quad (1.7)$$

It then follows by uniqueness of solutions to the Dirichlet problem (1.5) that

$$v^{(\ell)} := v_1^{(\ell)} = v_2^{(\ell)} \text{ in } \Omega, \quad \ell = 1, 2.$$



3. Second linearization.

For $j = 1, 2$, a straightforward computation shows that

$$\begin{cases} (\Delta + Q) w_j + 2a_j^{(2)} v^{(1)} v^{(2)} = 0 & \text{in } \Omega, \\ w_j = 0 & \text{on } \partial\Omega. \end{cases} \quad (1.8)$$

where $w_j = \partial_{\varepsilon_1 \varepsilon_2}^2 |_{\varepsilon=0} u_j$. Let $\mathbb{V}^{(\ell)}$ ($\ell = 1, 2$) be the solution of

$$\begin{cases} (\Delta + Q) \mathbb{V}^{(\ell)} = 0 & \text{in } \Omega, \\ \mathbb{V}^{(\ell)} = g_\ell & \text{on } \partial\Omega, \end{cases} \quad (1.9)$$

Multiply (1.8) by $\mathbb{V}^{(1)}$, then integration by parts gives rises to

$$\int_{\Omega} \left(a_1^{(2)} - a_2^{(2)} \right) v^{(1)} v^{(2)} \mathbb{V}^{(1)} dx = 0,$$

such that $\left(a_1^{(2)} - a_2^{(2)} \right) \mathbb{V}^{(1)} = 0$ in Ω . □

Semilinear reaction-diffusion equations

Consider the initial boundary value problem

$$\begin{cases} \rho(t, x) \partial_t u + \nabla \cdot (A(t, x) \nabla u) + b(t, x, u) = 0, & (t, x) \in (0, T) \times \Omega, \\ u(t, x) = f(t, x), & (t, x) \in (0, T) \times \partial\Omega, \\ u(0, x) = 0, & x \in \Omega. \end{cases}$$

Theorem (Kian-Liimatainen-L., 2023)

The same lateral DN map $\Lambda_{b_1}(f) = \Lambda_{b_2}(f)$, for all f , then there exists $\phi \in C^{1+\frac{\alpha}{2}, 2+\alpha}([0, T] \times \overline{\Omega})$ satisfying

$$\phi(0, x) = 0, \quad x \in \Omega, \quad \phi(t, x) = \partial_{\nu_a} \phi(t, x) = 0, \quad (t, x) \in \Sigma.$$

such that $b_1 = S_\phi b_2$, where S_ϕ is defined by

$$S_\phi b(t, x, \mu) = b(t, x, \mu + \phi(t, x)) + \rho(t, x) \partial_t \phi(t, x) + \nabla \cdot (A \nabla \phi).$$

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Fractional Laplacian

Let us start with some natural questions:

- What if the Laplacian $-\Delta$ is replaced by the **fractional Laplacian** $(-\Delta)^s$ for $0 < s < 1$? Can we consider the Calderón problem for the fractional Laplacian?

The fractional Laplacian is defined by

$$(-\Delta)^s u = \mathcal{F}^{-1} \{ |\xi|^{2s} \widehat{u}(\xi) \}, \text{ for all } u \in \mathcal{S}(\mathbb{R}^n).$$

The fractional Laplacian $(-\Delta)^s$ is **nonlocal**: It does not preserve the supports, and computing $(-\Delta)^s u(x)$ involves values of u away from x .

The exterior value problem: Fractional Schrödinger equation

Since $(-\Delta)^s$ is a **nonlocal** operator, the forward problem for the fractional Schrödinger equation is given by (for all dimension $n \in \mathbb{N}$)

$$\begin{cases} (-\Delta)^s u + qu = 0 & \text{in } \Omega, \\ u = f & \text{in } \Omega_e := \mathbb{R}^n \setminus \Omega. \end{cases}$$

The well-posedness of the above equation can be guaranteed by the Lax-Milgram. Hence, one can derive that the DN map of the fractional Schrödinger equation will be given by

$$\Lambda_q : f \mapsto (-\Delta)^s u|_{\Omega_e},$$

where u is the unique solution to the fractional Schrödinger equation.

The answer of the fractional Calderón problem is positive.

Theorem (Ghosh-Rüland-Salo-Uhlmann, single measurement)

Let $q_j \in C^0(\overline{\Omega})$, Λ_{q_j} be the DN maps of $(-\Delta)^s u + q_j u = 0$ in Ω , and $W_j \subset \mathbb{R}^n \setminus \overline{\Omega}$ be arbitrary open set, for $j = 1, 2$. Then $\Lambda_{q_1}(f)|_{W_2} = \Lambda_{q_2}(f)|_{W_2}$ for *one* $0 \neq f \in C_c^\infty(W_1)$. Then $q_1 = q_2$ in Ω .

Note that in the fractional case, the same DN map yields that

$$\int_{\Omega} (q_1 - q_2) u_1 u_2 dx = 0,$$

where u_1 and u_2 are the solutions in Ω with potentials q_1 and q_2 .

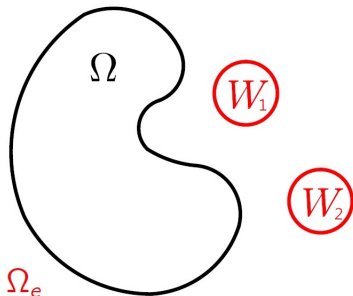
The result can be shown by the **unique continuation**:

Theorem (Unique continuation)

Let $\mathcal{O} \subset \mathbb{R}^n$ be an arbitrary open set. If $(-\Delta)^s u = 0$ in \mathcal{O} and $u = 0$ in a positive measurable subset of \mathcal{O} , one can conclude that $u \equiv 0$ in \mathbb{R}^n .

Main features

1. Partial data results for **arbitrary open sets** $W_1, W_2 \subset \Omega_e$ (W_1 and W_2 may not be disjoint)
2. The same method works for **any** dimensions $n \in \mathbb{N}$
3. New mechanism in solving nonlocal type inverse problems.



- Fractional Schrödinger equation (Ghosh-Salo-Uhlmann 2016).
- Nonlocal variable coefficients (Ghosh-L.Xiao 2017).
- Optimal stability (Rüland-Salo 2017).
- Nonlocal Schiffer (Cao-L.-Liu 2017).
- Monotonicity tests (Harrach-L. 2017, 2018).
- Single measurement and reconstruction (Ghosh-Rüland-Salo-Uhlmann 2018).
- Fractional conductivity (Covi 2018).
- Fractional Schrödinger equation with drift (Célic-L.-Rüland 2018).
- Lipschitz stability with finite dimension (Rüland-Sincich 2018).
- Fractional semilinear (Lai-L. 2018, 2020).
- Fractional heat equation (Lai-L.-Rüland 2019).
- Directionally antilocal principal symbols (Covi-García-Ferrero-Rüland 2021).
- Fractional wave equation (Kow-L.-Wang 2021).
- Nonlocal elliptic operators (Ghosh-Uhlmann 2021).
- Fractional anisotropic on closed manifolds (Feizmohammadi-Ghosh-Krupchyk-Uhlmann 2021).
- Inverse source and minimal number of measurements (Liu-L. 2022).
- Global uniqueness of conductivity (Covi-Railo-Zimmermann 2022).
- Counterexample constructions with disjoint measured sets (Railo-Zimmermann 2022).
- Low regularity for γ (Railo-Zimmermann 2022).
- Nonlocal parabolic equation (Banerjee-Krishnan-Senapati 2022).
- Logarithmic stability (Covi-Railo-Tyni-Zimmermann 2022).
- Nonlocal parabolic operators (Lin-L.-Uhlmann 2022).
- Fractional elasticity (Covi-de Hoop-Salo 2022).
- Fractional p -Laplacian (Kar-L.-Zimmermann 2022).

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Fractional elliptic equations

Consider the fractional equation

$$\begin{cases} (-\Delta)^s u(x) + a(x, u) = F & \text{in } \Omega, \\ u = f & \text{in } \Omega_e, \end{cases} \quad (3.1)$$

where $0 < s < 1$. We are interested to determine

$$a(x, u) := \sum_{k=1}^N a^{(k)}(x)[u(x)]^k,$$

for $N \in \mathbb{N} \cup \{0\} \cup \{\infty\}$, and F .

When $a(x, u)$ is **nonlinear**, we need to assume the condition

$$a^{(1)}(x) = \partial_u a(x, 0) \geq 0 \text{ for } x \in \Omega,$$

in order to prove the local well-posedness of (3.1).

The case $s = 1$ and $a^{(k)} = 0$

Consider

$$\begin{cases} -\Delta u_j = F_j & \text{in } \Omega, \\ u_j = f & \text{on } \partial\Omega, \end{cases} \quad (3.2)$$

for $j = 1, 2$. In fact, to find the obstruction for the unique determination problem, let $\phi \in C_c^2(\Omega)$ be an arbitrary function, then one has $\phi = \partial_\nu \phi = 0$ on $\partial\Omega$. Let $(u_j|_{\partial\Omega}, \partial_\nu u_j|_{\partial\Omega})$ be the Cauchy data of (3.2), even if

$$(u_1|_{\partial\Omega}, \partial_\nu u_1|_{\partial\Omega}) = (u_2|_{\partial\Omega}, \partial_\nu u_2|_{\partial\Omega}),$$

but we can only prove the **gauge invariance** that $F_2 = F_1 - \Delta\phi$, and $\Delta\phi$ can be arbitrary. Therefore, the unique determination is not possible for the unknown sources in general.

Nonlocal Cauchy data

Let W_1 and W_2 be two arbitrary nonempty open subsets in Ω_e . It is always assumed that $\text{supp}(f) \subset W_1$, and moreover $f|_{W_1} \in C^{2,s}(\overline{W_1})$. With the well-posedness at hand, we introduce the following exterior **nonlocal partial Cauchy data** set:

$$\begin{aligned}\mathcal{C}_{a,F}^{W_1,W_2}(f) &:= (u|_{W_1}, (-\Delta)^s u|_{W_2}) \\ &= (f|_{W_1}, (-\Delta)^s u|_{W_2})\end{aligned}\tag{3.3}$$

where $u \in C^s(\mathbb{R}^n)$ is the unique solution to (3.1).

- The following results are holds at least when $\Omega \subset \mathbb{R}^n$ is a $C^{1,1}$ bounded domain and for any dimension $n \in \mathbb{N}$.

Minimal number of measurements

Theorem (L.-Liu, 2022)

Let $W_1, W_2 \subset \Omega_e$ be two arbitrary nonempty open subsets, and consider

$$\begin{cases} (-\Delta)^s u_j + a_j(x, u_j) = F_j & \text{in } \Omega, \\ u_j = f & \text{in } \Omega_e, \end{cases} \quad (3.4)$$

where $a_j(x, u_j) = \sum_{k=1}^N a_j^{(k)} [u_j(x)]^k$, $j = 1, 2$ and a *finite* $N \in \mathbb{N}$. Assuming the well-posedness of (3.4), if

$$\mathcal{C}_{a_1, F_1}^{W_1, W_2}(f_k) = \mathcal{C}_{a_2, F_2}^{W_1, W_2}(f_k), \quad k = 0, 1, \dots, N, \quad (3.5)$$

where $f_k \not\equiv f_l$, $0 \leq k, l \leq N$ and $k \neq l$, then one has

$$a_1^{(k)}(x) = a_2^{(k)}(x) \text{ in } \Omega, \quad k = 1, 2, \dots, N, \quad \text{and} \quad F_1 = F_2 \text{ in } \Omega.^a$$

^aThe number of unknowns equal to the number of measurements

The Proof

We want to show that

$$a_1^{(k)} = a_2^{(k)} \text{ in } \Omega, \text{ for } k = 1, 2, \dots, N. \quad (3.6)$$

Let $f_0 = 0, f_1, \dots, f_N \in Y_\delta$, which are mutually different, and consider $u_j^{(\ell)}$ to be the solutions of

$$\begin{cases} (-\Delta)^s u_j^{(\ell)} + a_j(x, u_j^{(\ell)}) = F_j & \text{in } \Omega, \\ u_j^{(\ell)} = f_\ell & \text{in } \Omega_e, \end{cases} \quad (3.7)$$

for $\ell = 0, 1, \dots, N$ and $j = 1, 2$.

By the strong uniqueness for the fractional Laplacian,

$$u^{(\ell)} := u_1^{(\ell)} = u_2^{(\ell)} \text{ in } \mathbb{R}^n, \quad \text{for } \ell = 0, 1, \dots, N. \quad (3.8)$$

Moreover, via (3.7) and (3.8), it is not hard to derive

$$\sum_{k=1}^N a_j^{(k)} \left(u^{(\ell)} - u^{(0)} \right)^k = 0 \text{ in } \Omega,$$

for $j = 1, 2$, so that

$$\sum_{k=1}^N \left(a_1^{(k)} - a_2^{(k)} \right) \left(u^{(\ell)} - u^{(0)} \right)^k = 0 \text{ in } \Omega, \quad (3.9)$$

for all $\ell = 0, 1, \dots, N$.

Rewrite (3.9) as $UA = 0$ in Ω , where U is an $N \times N$ matrix

$$U := \begin{pmatrix} u^{(1)} - u^{(0)} & (u^{(1)})^2 - (u^{(0)})^2 & \dots & (u^{(1)})^N - (u^{(0)})^N \\ u^{(2)} - u^{(0)} & (u^{(2)})^2 - (u^{(0)})^2 & \dots & (u^{(2)})^N - (u^{(0)})^N \\ \vdots & \vdots & \ddots & \vdots \\ u^{(N)} - u^{(0)} & (u^{(N)})^2 - (u^{(0)})^2 & \dots & (u^{(N)})^N - (u^{(0)})^N \end{pmatrix} \quad (3.10)$$

and A is an N -column vector

$$A := \begin{pmatrix} a_1^{(1)} - a_2^{(1)} \\ a_1^{(2)} - a_2^{(2)} \\ \vdots \\ a_1^{(N)} - a_2^{(N)} \end{pmatrix}. \quad (3.11)$$

It suffices to show that the matrix U in (3.10) is **non-singular** a.e. in Ω .

Via direct computations, we have

$$\det U = \det \begin{pmatrix} u^{(1)} - u^{(0)} & (u^{(1)})^2 - (u^{(0)})^2 & \dots & (u^{(1)})^N - (u^{(0)})^N \\ u^{(2)} - u^{(0)} & (u^{(2)})^2 - (u^{(0)})^2 & \dots & (u^{(2)})^N - (u^{(0)})^N \\ \vdots & \vdots & \ddots & \vdots \\ u^{(N)} - u^{(0)} & (u^{(N)})^2 - (u^{(0)})^2 & \dots & (u^{(N)})^N - (u^{(0)})^N \end{pmatrix}$$

$$= \det \begin{pmatrix} 1 & u^{(0)} & (u^{(0)})^2 & \dots & (u^{(0)})^N \\ 0 & u^{(1)} - u^{(0)} & (u^{(1)})^2 - (u^{(0)})^2 & \dots & (u^{(1)})^N - (u^{(0)})^N \\ 0 & u^{(2)} - u^{(0)} & (u^{(2)})^2 - (u^{(0)})^2 & \dots & (u^{(2)})^N - (u^{(0)})^N \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & u^{(N)} - u^{(0)} & (u^{(N)})^2 - (u^{(0)})^2 & \dots & (u^{(N)})^N - (u^{(0)})^N \end{pmatrix}$$

Hence,

$$\det U = \det \begin{pmatrix} 1 & u^{(0)} & (u^{(0)})^2 & \dots & (u^{(0)})^N \\ 1 & u^{(1)} & (u^{(1)})^2 & \dots & (u^{(1)})^N \\ 1 & u^{(2)} & (u^{(2)})^2 & \dots & (u^{(2)})^N \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & u^{(N)} & (u^{(N)})^2 & \dots & (u^{(N)})^N \end{pmatrix},$$

which is the [Vandermonde matrix](#) in the last identity and

$$\det U = \prod_{1 \leq \ell < m \leq N} (u^{(m)} - u^{(\ell)}) \neq 0 \text{ a.e. in } \Omega.$$

Therefore, one can conclude that the vector A in (3.11) must be zero a.e. in Ω . Since each $a_j^{(k)} \in C^s(\overline{\Omega})$, for $j = 1, 2$, $k = 1, 2, \dots, N$, the claim (3.6) must hold. Finally, by using the equation (3.7), we can summarize that $F_1 = F_2$ in Ω as well.

Conclusions

- One can challenge open/unsolved inverse problems under fractional settings.
- Nonlocality is **beneficial** in solving related inverse problems.
- Some features of inverse problems are given:
 - ▶ **Local.** Recover coefficients then solutions.
 - ▶ **Nonlocal.** Recover solutions then coefficients.
 - ★ **No** linearization techniques are involved.
- As $s = 1$, $k = 2$, we² can only prove that there exists $\phi \in C_0^2(\Omega)$ such that there is a gauge invariance.
- On other hand, the uniqueness result for source holds for a nonlocal model³.
- Gauge symmetry and breaking for parabolic equations are investigated.⁴
- Future study: Optimal condition to break the gauge.

²Liimatainen-L., 2022

³L.-Liu, 2022

⁴Kian-Liimatainen-L., in 2 weeks

Thank you for your attention !